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Isolation Spectroscopic Characterization of 3-Methyl-2-Nitroanisole by Density Functional Method

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ABSTRACT

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A combined experimental and theoretical studies were conducted on the molecular structure and vibrational, spectra of 3-methyl-2-nitro anisole (MNA). The FT-IR and FT-Raman spectra of (MNA) were recorded in the solid phase. The molecular geometry and vibrational frequencies of MNA in the ground state have been calculated by using the abinitio HF (Hartree-Fock) and density functional methods (B3LYP) invoking 6-31+G (d,p) basis set. The optimized geometric bond lengths and bond angles obtained by HF method shows best agreement with the experimental values. Comparison of the observed fundamental vibrational frequencies of MNA with calculated results by HF and density functional methods indicates that B3LYP is superior to the scaled HF approach for molecular vibrational problems. The difference between the observed and scaled wave number values of most of the fundamental is very small. The thermodynamic functions and atomic change of the title compound was also performed at HF/B3LYP/6-31+G(d,p) level of theories. A detailed interpretation of the NBO, NMR spectra of MNA was also reported. The thermodynamic function of the title compound was also performed at HF/6-31+G (d,p) and B3LYP/6-31+G (d,p) level of theories. Natural bond orbital analysis has been carried out to explain the change transfer or delocalization of change due to the intra-molecular interactions. Energy of the highest occupied molecular (HOMO) orbital and lowest unoccupied (LUMO) molecular orbital have been predicted.

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1. Introduction

Anisole is a crystalline powder that is soluble in ether and alcohol insoluble in water; boiling point 155°C. Anisole is more electron rich than benzene because of the resonance effect of the methoxy group¹. Anisole reacts with electrophilic aromatic substitution reaction quickly than benzene. Anisole has the two effects, i.e. back-donation and conjugation, can be competitive and their relative weight determining the electronic properties of the molecule can be conformation independent². Anisole is used in perfumery, chemical syntheses and an insect pheromone. Particularly anisole is a bio-agent and it is mainly used to synthesize amino anisole, dyes and medicine. The application of anisole as a detector for middle infrared interferometry ³ has been confirmed. Since the anisole absorbs IR radiation heat it is possible to estimate the IR intensity distribution on the anisole from the diffraction pattern made by visible larger light that is transmitted through it. It is used to measure the mono phenolase activity of polyphenol oxidize from fruits and vegetables⁴. Other antioxidants for food are phosphoric acid, citric acid, gallic acid, ascorbic acid and their esters which form complexes with the pro-oxidative metal traces. Antimicrobial process is also important in preserving foods. MNA are used as an antioxidant in plastics, elastomers and petroleum (lubes, greases and waxes), practically bigger market size than food field. MNA is used as a stabilizer to inhibit the autopolymerization of organic peroxides. More recently, the experimental vibrational spectra of p-nitroanisole, 2-

nitroanisole, thioanisole, and 3-nitroanisole⁵ have been investigated in comparison with DFT/B3LYP level calculations.In the case of IR laser interferometry, interference fringe patterns that are created by the IR laser on the anisole are observed as the refractive index distribution; hence the anisole functions as a phase grating for visible light. It is, therefore, conceivable to estimate the IR intensity distribution on the anisole by analyzing the diffraction pattern made by a visible laser, which is transmitted through the anisole. Also, since a two dimensional device using the anisole does not require matrix structure, it is expected that the measurement system must have high spatial resolution, equivalent to that of existing IR cameras. Owing to these applications and the reliable properties of anisole, a complete vibrational study on 3-methyl-2-nitroanisole (MNA) has been undertaken. The vibrational analyses of MNA using the SQM force field method based on DFT calculation have been presented. The calculated infrared and Raman spectra of MNA are also simulated utilizing the computed dipole derivatives for IR intensities and polarizability derivatives for Raman activities.

2. Experimental Analysis

The chemicals required were obtained from Sigma Aldrich chemical suppliers and are of analar grade. FTIR spectra (KBr pellets) were recorded on 8101 Schimadzu FTIR spectrophotometer recorded in the region 4000 - 400 cm⁻¹. FT-Raman spectra have been recorded in the region 3500 - 50 cm⁻¹ on a Perkin-Elmer spectrometer.

3. Computational Methodology

The entire calculations are performed at Hartree-Fock (HF) and density functional (DFT) levels using GAUSSIAN 09W⁶ program package, invoking gradient geometry optimization ⁷. Initial geometry generated from the standard geometrical parameters was minimized without any constraint on the potential energy surface at HF level adopting the standard 6-311++G(d,p) basis set. This geometry was then reoptimized again at DFT level employing the Becke 3LYP keyword, which invokes Becke's three-parameter hybrid method⁸ using the correlation function of Lee *et al.*⁹; implemented with the same basis set. The optimized structural parameters are used for the vibrational frequency calculations at DFT level to characterize all the stationary points as minima. The multiple scaling of the force constants are performed according to SOM procedure ^{10,11} using selective scaling in the natural internal coordinate representation ^{12,13} The transformation of force field; subsequent normal coordinate analysis and calculation of the TED are done on a PC with MOLVIB program (version V7.0 – G77) written by Sundius ¹⁴⁻¹⁶. By the use of GAUSSVIEW molecular visualization program ¹⁷ along with available related molecules; the vibrational frequency assignments are made by their TED with a high degree of confidence. The TED elements provide a measure of each internal co-ordinate's contributions to the normal coordinate.

3.1. Prediction of Raman intensities

The Raman activities (S_i) calculated with the GAUSSIAN 09W program are subsequently converted to relative Raman intensities (I_i) using the following relationship derived from the basic theory of Raman scattering ¹⁸⁻²⁰

$$\mathbf{I}_{i} = \frac{f\left(\mathbf{v}_{o} - \mathbf{v}_{i}\right)^{4} \mathbf{S}_{i}}{\mathbf{v}_{i} \left[1 - \exp\left(-hcv_{i}/kT\right)\right]} \qquad \dots (1)$$

where v_0 is the exciting frequency in cm⁻¹, v_i the vibrational wave number of the *i*th normal mode, *h*, *c* and *k* are the fundamental constants and *f* is a suitably chosen common normalisation factor for all the peak intensities.

4. Results and Discussion 4.1. Molecular geometry

The molecular structure of MNA belongs to C_1 point group symmetry. For C_1 symmetry there would not be any relevant distribution. The molecule consists of 21 atoms and expected to have 57 normal modes of vibrations of the same a species under C_1 symmetry. These modes are found to be IR and Raman active suggesting that the molecule possesses a

and Raman active suggesting that the molecule possesses a non-centrosymmetric structure, which recommends the MNA for non-linear optical applications. The optimized molecular structure of MNA is shown in Fig. 1, respectively.



Fig 1. Molecular structure of 3-methyl-2-nitro anisole.

4.2. PES scan studies

The potential energy surface (PES) scan with the B3LYP/6-311++G(d,p) level of theoretical approximations was performed for MNA is shown in Fig. 2, respectively. The dihedral angle C1-C2-N12-O14 for MNA is also relevant coordinate for conformational flexibility within the molecule. During the calculation, all the geometrical parameters are simultaneously relaxed while the C1-C2-N12-O14 torsional angle are varied in steps of 10°, 20°, 30°,..., 360°. For this rotation minimum energy curves have been obtained at 0° . 110° and 250° as shown in Figs. 3 and 4 clearly demonstrates that 0° corresponds to the global minimum energy is -590.54807307 Hartrees for MNA, respectively and the minimum energy for the rotation at 110° and 250° are -590.54780773, -590.54784943 Hartrees for MNA, respectively. The maximum energy is obtained at 180° for both the molecules.



Fig 2. PES scan of 3-methyl-2-nitro anisole. 4.3. Structural properties

The optimization geometrical parameters of MNA obtained by the ab initio HF and DFT/B3LYP methods with 6-311++G(d,p) as basis set are listed in Table 1. The computed bond length and bond angles are compared with Xray diffraction data of similar compound ²¹. From Table 1, it can be seen that there are some deviations in the computed geometrical parameters from those reported in the single crystal XRD data, and these differences are probably due to the intermolecular interactions in the crystalline state. Comparing bond angles and bond lengths of B3LYP method with HF method, various bond lengths are found to be almost same at HF/6-311++G(d,p) and B3LYP/6-311++G(d,p) levels. However, the B3LYP/6-311++G(d,p) level of theory, in general slightly over estimates bond lengths but it yields bond angles in excellent agreement with the HF method. For the investigated molecule, the benzene ring appears a little distorted with larger C4-C5, C1-C2 bond length and shorter C1-C2, C4-C5 bond length and angles slightly out of the regular hexagonal structure for MNA. These distortions are explained in terms of the change in hybridisation affected by the substituent at the carbon site to which it is appended. The MNA of C-C bond lengths adjacent to the C3-O15, C8-O7 bonds are increases and the angles C2-C1-C6, C2-C3-C4, C1-C6-C5 are smaller than typical hexagonal angle of 120°. This is because of the effect of substitution of O-CH₃ groups attached to the C6 and C1 of the benzene ring for MNA. The variation in torsional angles C6- $C1-O7-C8 = 77.01^{\circ}$ (HF), -2.15° (B3LYP) is due to charge delocalization for MNA. The calculated geometric parameters can be used as foundation to calculate the other parameters for the molecule.

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| Table 1. Optimized geometrical | parameters of 3-methyl-2-nitroanisole | e by using B3LYP/6-311+G(d | l.p) and B3LYP/ 6-311++G(d. | p) methods and basis set calculations. |
|--------------------------------|---------------------------------------|----------------------------|-----------------------------|--|
| | | | \mathbf{T} | |

| Bond Length | Value (Å) | | Exp | Bond Angle | Value (°) | | Value (°) Exp | | Value (°) | |
|-------------|---------------|--------------|-------|------------|---------------|----------------|---------------|---------------|---------------|---------------|
| | HF/ | B3LYP/ | | | HF/ | B3LYP/ | | | HF/ | B3LYP/ |
| | 6-311++G(d,p) | 6-11++G(d,p) | | | 6-311++G(d,p) |)6-311++G(d,p) | | | 6-311++G(d,p) | 6-311++ (d,p) |
| C1-C2 | 1.3862 | 1.4028 | | C2-C1-C6 | 118.5865 | 118.1802 | | C6-C1-C2-C3 | -0.0735 | -0.473 |
| C1-C6 | 1.3866 | 1.3966 | | C2-C1-O7 | 119.7595 | 116.589 | | C6-C1-C2-N12 | 179.9527 | 179.2257 |
| C1-07 | 1.3485 | 1.3527 | | C6-C1-O7 | 121.5198 | 125.2154 | | O7-C1-C2-C3 | -175.9267 | 178.1723 |
| C2-C3 | 1.3903 | 1.3936 | | C1-C2-C3 | 123.3382 | 123.3088 | 117.82 | O7-C1-C2-C12 | 4.0995 | -2.1291 |
| C2-N12 | 1.4641 | 1.4807 | 1.491 | C1-C2-N12 | 117.5613 | 117.7382 | 118.6 | C2-C1-C6-C5 | 1.0595 | -0.3454 |
| C3-C4 | 1.3892 | 1.3999 | | C3-C2-N12 | 119.1005 | 118.9523 | | C2-C1-C6-H21 | -177.959 | 179.6272 |
| C3-C15 | 1.511 | 1.508 | | C2-C3-C4 | 116.6463 | 117.0158 | | O7-C1-C6-C5 | 176.8364 | -178.8625 |
| C4-C5 | 1.3872 | 1.3898 | | C2-C3-C15 | 122.4958 | 121.8525 | | O7-C1-C6-H21 | -2.1821 | 1.1101 |
| C4-H19 | 1.0748 | 1.0835 | 1.090 | C4-C3-C15 | 120.8559 | 121.13 | | C2-C1-O7-C8 | -107.2579 | 179.3024 |
| C5-C6 | 1.3829 | 1.3934 | | C3-C4-C5 | 121.0881 | 120.7999 | | C6-C1-O7-C8 | 77.0138 | -2.1593 |
| С5-Н20 | 1.0753 | 1.084 | 1.090 | C3-C4-H19 | 119.1763 | 119.0847 | | C1-C2-C3-C4 | -1.0861 | 1.1012 |
| C6-H21 | 1.0743 | 1.0811 | 1.091 | C5-C4-H19 | 119.7346 | 120.1154 | | C1-C2-C3-C15 | 179.4303 | -179.3608 |
| O7-C8 | 1.4172 | 1.4246 | | C4-C5-C6 | 120.8808 | 121.2027 | | N12-C2-C3-C4 | 178.8874 | -178.594 |
| С8-Н9 | 1.0823 | 1.0946 | 1.090 | C4-C5-H20 | 119.5668 | 119.7141 | | N12-C2-C3-C15 | -0.5962 | 0.944 |
| C8-H10 | 1.0794 | 1.0879 | 1.090 | С6-С5-Н20 | 119.5492 | 119.0831 | | C1-C2-N12-O13 | -177.2474 | 113.0626 |
| C8-H11 | 1.0845 | 1.0944 | 1.090 | C1-C6-C5 | 119.4417 | 119.483 | | C1-C2-N12-O14 | 63.9021 | -68.2033 |
| N12-O13 | 2.6681 | 1.224 | 1.227 | C1-C6-H21 | 119.2362 | 120.5328 | | C3-C2-N12-O13 | 62.7776 | -67.2252 |
| N12-O14 | 1.1924 | 1.2205 | 1.228 | C5-C6-H21 | 121.3146 | 119.9842 | | C3-C2-N12-O14 | -116.0729 | 111.5089 |
| C15-H16 | 2.7286 | 1.0904 | | C1-07-C8 | 117.0894 | 118.9148 | 117.64 | C2-C3-C4-C5 | 1.2815 | -0.9304 |
| C15-H17 | 1.0824 | 1.0934 | | O7-C8-H9 | 110.3512 | 111.2231 | 117.80 | C2-C3-C4-H19 | -179.0757 | 179.0448 |
| C15-H18 | 1.0832 | 1.0926 | | O7-C8-H10 | 106.0602 | 105.5137 | 117.80 | C15-C3-C4-C5 | -179.2259 | 179.5281 |
| | | | | O7-C8-H11 | 110.7278 | 111.129 | 117.80 | С15-С3-С4-Н19 | 0.4169 | -0.4968 |
| | | | | H9-C8-H10 | 109.8492 | 109.5129 | 124.31 | C2-C3-C15-H16 | 166.7484 | -170.8561 |
| | | | | H9-C8-H11 | 110.1785 | 109.8752 | | C2-C3-C15-H17 | -73.1455 | -50.6739 |
| | | | | H10-C9-H11 | 109.5918 | 109.4847 | | C2-C3-C15-H18 | 46.8301 | 68.859 |
| | | | | C2-N12-O13 | 112.7465 | 117.1751 | | C4-C3-C15-H16 | -12.7139 | 8.6631 |
| | | | | C2-N12-O14 | 117.4204 | 117.4914 | | C4-C3-C15-H17 | 107.3922 | 128.8453 |

^a For numbering of atoms refer Fig. 2. ^a Experimental values are taken from Ref. [21].

5. vibrational spectra

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Detailed description of vibrational modes can be given by means of normal coordinate analysis. For this purpose, the full set of 71 standard internal coordinates (containing 14 redundancies) for MNA are presented in Table 2, respectively. From these, a non-redundant set of local symmetry coordinates are constructed by suitable linear combinations of internal coordinates following the recommendations of Fogarasi et al. and they are summarized in Table 3 for MNA, respectively. The detailed vibrational assignment of fundamental modes of MNA along with the calculated IR, frequencies and normal mode descriptions Raman (characterized by TED) are reported in Table 4. For visual comparison, the observed and calculated FTIR and FT-Raman spectra of MNA at HF and B3LYP levels using 6-311++G(d,p) basis set are shown in Figs. 3 - 4. The vibrational analysis obtained for MNA with the unscaled HF and B3LYP/6-311++G(d,p) force field are generally somewhat greater than the experimental values due to neglect of anharmonicity in real system. These discrepancies can be corrected either by computing anharmonic corrections explicitly or by introducing a scaled field or directly scaling the calculated wave numbers with proper factor ²². A tentative assignment is often made on the basis of the unscaled frequencies by assuming the observed frequencies so that they are in the same order as the calculated ones. Then, for an easier comparison to the observed values, the calculated frequencies are scaled by the scale to less than 1, to minimize the overall deviation. The results indicate that the B3LYP calculations approximate the observed fundamental frequencies much better than the HF results. Inclusion of

electron correlation in density functional theory to a certain extent makes the frequency values smaller in comparison with HF frequency. Also, it should be noted that the experimental results belong to solid phase and theoretical calculations belong to gaseous phase. A better agreement between the computed and experimental frequencies can be obtained by using different scale factors for different regions of vibrations. For that purpose, we have utilized different scaling factors for all fundamental modes except the torsional mode to obtain the scaled frequencies of the compound. The resultant scaled frequencies are also listed in Table 4.



Fig 3. Comparison of observed and calculated infrared spectra of 3-methyl-2-nitro anisole.

(a) Observed

(b) HF/6-311++G(d,p) (c) B3LYP/6-311++G(d,p)

| No. (i) | Symbol | Туре | Definition ^a |
|----------------------|------------------|--------------------------|--|
| Stretching | - | | |
| 1-3 | r _i | СН | C4-H19, C5-H20, C6-H21 |
| 4-9 | η_i | CH ₃ (methyl) | C8-H9, C8-H10, C8-H11, C15-H16, C15-H17, C15-H18 |
| 10-11 | R _i | СО | C1-07, C8-07 |
| 12-18 | q_i | CC | C1-C2, C2-C3, C3-C4, C4-C5, C5-C6, |
| | _ | | C6-C1, C3-C15 |
| 19 | Si | CN | C2-N12 |
| 20-21 | Xi | NO | N12-O13, N12-O14 |
| In-plane bending | | | |
| 22-27 | α_i | Ring | C1-C2-C3, C2-C3-C4, C3-C4-C5, C4-C5-C6, |
| | | | C5-C6-C1, C6-C1-C2 |
| 28-33 | β_i | ССН | C1-C6-H21, C5-C6-H21, C6-C5-H20, |
| | | | C4-C5-H20, C5-C4-H19, C3-C4-H19 |
| 34-35 | π_{i} | CCC | C2-C3-C15, C4-C3-C15 |
| 36-37 | θ_i | CCO | C6-C1-O7, C2-C1-O7 |
| 38 | θ_i | COC | C1-07-C8 |
| 39-41 | υ_i | CCH(methyl) | С3-С15-Н16, С3-С15-Н17, С3-С15-Н18 |
| 42-44 | β_i | НСН | H16-C15-H17, H17-C15-H18, H16-C15-H18 |
| 45-47 | υ_i | OCH | O7-C8-H9, O7-C8-H10, O7-C8-H11 |
| 48-50 | β_i | НСН | H9-C8-H10, H10-C8-H11, H9-C8-H11 |
| 51-52 | \mathbf{f}_{i} | CCN | C3-C2-N12, C1-C2-N12 |
| 53-54 | ψ_i | CNO | C2-N12-O13, C2-N12-O14 |
| 55 | ρ | O-N-O | O13-N12-O14 |
| Out-of-plane bending | | | |
| 56-58 | δ_i | СН | H19-C5-C4-C3, H20-C5-C6-C4, H21-C6-C1-C5 |
| 59 | Ψ_i | OC | 07-C1-C2-C6 |
| 60 | λ_i | OC | C8-O7-C1-(C6-C2) |
| 61 | δ_i | CC | C15-C3-C2-C4 |
| 62 | χi | CN | N12-C2-C3-C1 |
| Torsion | | | |
| 63-68 | $	au_{i}$ | tRing | C1-C2-C3-C4, C2-C3-C4-C5, C3-C4-C5-C6, |
| | | | C4-C5-C6-C1, C5-C6-C1-C2, C6-C1-C2-C3 |
| 69 | τ_{i} | tC-CH ₃ | (C2-C4)-C3-C15-(H16,H17,H18) |
| 70 | τ_{i} | tO-CH ₃ | C1-O7-C8-(H9,H10,H11) |
| 71 | τ_i | tC-NO ₂ | (C1-C2)-C2-N12-(O13,O14) |

^aFor numbering of atoms refer Fig. 2.

C−*H vibrations* □ The hetero aromatic structure shows the presence of C−H stretching vibration in the region 3100 – 3000 cm⁻¹ which is the characteristic region for the identification of such C−H stretching vibrations. These vibrations are not found to be affected due to the nature and position of the substituents. In the present investigation, the C−H vibrations are observed at 3101, 3082, 3030 cm⁻¹ in FT-Raman spectrum for MNA and the corresponding force constant contribute to the TED for MNA. The C-H in-planebending vibrations usually occur in the region 1390 - 990 cm⁻¹ and are very useful for characterization purposes. The C−H out-of-plane bending vibrations occur in the region 900 - 675 cm⁻¹ and these bands are highly informative ²². Accordingly, the 998, 976 cm⁻¹ in the FTIR spectrum and 1001 cm⁻¹ in Raman spectrum for MNA have been assigned to, C–H inplane-bending vibrations. The C–H out-of-plane mode is observed at 776, 745 cm⁻¹ in the FTIR and at 752 cm⁻¹ in Raman spectrum for MNA.

The observed C–H out-of-plane bending modes show consistent agreement with the computed B3LYP and HF results. *C–C vibrations* \Box The C–C hetero aromatic stretching vibrations generally occur in the region 1650 - 1400 cm⁻¹²². With heavy substituents, the bonds tend to shift to somewhat lower wavenumbers and greater the number of substituents on the ring, broader the absorption regions. As predicted in the earlier references, in the present investigation, the C–C stretching vibrations for MNA, the FTIR peaks observed at 1586, 1522, 1374 and 1310 cm⁻¹ and in FT-

| No | Туре | Definition ^a |
|-------|-----------------------|--|
| 1-3 | СН | r_1, r_2, r_3 |
| 4,5 | CH ₃ ss | $\eta_4 + \eta_5 + \eta_6 / \sqrt{3}$, $\eta_7 + \eta_8 + \eta_9 / \sqrt{3}$ |
| 6,7 | CH ₃ ips | $2\eta_4 + \eta_5 + \eta_6 / \sqrt{6}$, $2\eta_7 + \eta_8 + \eta_9 / \sqrt{6}$ |
| 8,9 | CH ₃ ops | $\eta_{5} - \eta_{6} / \sqrt{6}$, $\eta_{8} - \eta_{9} / \sqrt{6}$ |
| 10,11 | СО | R ₁₀ , R ₁₁ |
| 12-18 | CC | $q_{12}, q_{13}, q_{14}, q_{15}, q_{16}, q_{17}, q_{18}$ |
| 19 | CN | S_{20} |
| 20 | NO ₂ ss | $X_{20} - X_{21} / \sqrt{2}$ |
| 21 | NO ₂ ass | $X_{20} + X_{21}/\sqrt{2}$ |
| 22 | Rtrigd | $(\alpha_{22} - \alpha_{23} + \alpha_{24} - \alpha_{25} + \alpha_{26} - \alpha_{27}) / \sqrt{6}$ |
| 23 | Rsymd | $(-\alpha_{22} - \alpha_{23} + 2\alpha_{24} - \alpha_{25} - \alpha_{26} + 2\alpha_{27}) / \sqrt{12}$ |
| 24 | Rasymd | $(\alpha_{22} - \alpha_{23} + \alpha_{25} - \alpha_{26}) / \sqrt{2}$ |
| 25-27 | bCH | β_{28} - β_{29} / $\sqrt{2}$, β_{30} - β_{31} / $\sqrt{2}$, β_{32} - β_{33} / $\sqrt{2}$ |
| 28 | bCC | $\pi_{34} - \pi_{35}/\sqrt{2}$ |
| 29 | bCO | $\theta_{36} - \theta_{37}/\sqrt{2}$ |
| 30 | bOC | θ_{38} |
| 31-32 | CH ₃ sb | $(-\upsilon_{39} - \upsilon_{40} - \upsilon_{41} + \beta_{42} + \beta_{43} + \beta_{44})/\sqrt{2}$ |
| | | $(-\upsilon_{45} - \upsilon_{46} - \upsilon_{47} + \beta_{48} + \beta_{49} + \beta_{50})/\sqrt{2}$ |
| 33-34 | CH ₃ ipb | $(2\beta_{44}$ - β_{43} - β_{42} / $\sqrt{6}$, (2 β_{50} - β_{49} - β_{48})/ $\sqrt{6}$ |
| 35-36 | CH ₃ opb | $(\beta_{42} - \beta_{43} / \sqrt{6}, (\beta_{48} - \beta_{49} / \sqrt{6})$ |
| 37-38 | CH ₃ ipr | $(2\upsilon_{39} + \upsilon_{40} - \upsilon_{41})/\sqrt{6}$, $(2\upsilon_{45} + \upsilon_{46} - \upsilon_{47})/\sqrt{6}$ |
| 39-40 | CH ₃ opr | $(\upsilon_{40} - \upsilon_{41})/\sqrt{2}$, $(\upsilon_{46} - \upsilon_{47})/\sqrt{2}$ |
| 41 | £CN | $(f_{51} - f_{52})/\sqrt{2}$ |
| 42 | NO ₂ rock | $\psi_{53} - \psi_{54}/\sqrt{2}$ |
| 43 | NO ₂ twist | $\psi_{53} - \psi_{54}/\sqrt{2}$ |
| 44 | NO ₂ sciss | $2\rho_{55} - \psi_{54} - \psi_{53}/\sqrt{2}$ |
| 45-47 | δCH | $\delta_{56}, \delta_{57}, \delta_{58}$ |
| 48 | ψCO | Ψ59 |
| 49 | λΟC | λ_{60} |
| 50 | λCC | δ_{61} |
| 51 | χCN | χ ₆₂ |
| 52 | tKtrigd | $(\tau_{63} - \tau_{64} + \tau_{65} - \tau_{66} + \tau_{67} - \tau_{68}) / \sqrt{6}$ |
| 53 | tRsymd | $(\tau_{63} - \tau_{65} + \tau_{67} - \tau_{68}) / \sqrt{2}$ |
| 54 | tRasymd | $(-\tau_{63} + 2\tau_{64} - \tau_{65} - \tau_{66} + 2\tau_{67} - \tau_{68}) / \sqrt{12}$ |
| 55 | tCH ₃ | τ ₆₉ |
| 56 | tOCH ₃ | τ ₇₀ |
| 57 | NO_2wag | τ_{71} |

 Table 3. Definition of local symmetry coordinates of 3-methyl-2-nitroanisole.

^a The internal coordinates used here are defined in Table 4.

Raman the peaks are observed at 1559, 1538, 1489, 1374 and 1311 cm^{-1} are assigned for C–C stretching vibrations are confirmed by their TED values. Most of the ring vibrational modes are affected by the substitutions in the hetero aromatic ring of MNA. The bands observed at 728 cm⁻¹ in the FTIR for MNA have been designated to ring in-plane bending modes by careful consideration of their quantitative descriptions. The ring out-of-plane bending modes of MNA are also listed in Table 4 respectively.



Fig 4.Comparison of observed and calculated Raman spectra of 3-methyl-2-nitro anisole.

(a) Observed

(b) HF/6-311++G(d,p)

(c) B3LYP/6-311++G(d,p)

C–O vibrations \Box The interaction of the carbonyl group with a hydrogen donor group does not produce drastic and characteristic changes in the frequency of the C=O stretch as does by O–H stretch. A great deal of structural information can be derived from the exact position of the carbonyl stretching absorption peak. Susi and Ard ²² identified the C=O stretching mode at 1645 and 1614 cm⁻¹. On referring to the above findings and on the basis of the results of the normal coordinate analysis, the present investigation, the C–O

stretching vibrations have been found at 1189, 1178 cm⁻¹ in IR for MNA are assigned for C–O stretching vibrations are confirmed by their TED values. The C–O in-plane and out-of-plane bending vibrations level also have been identified and presented in Table 4 respectively for MNA.

 CH_3 group vibrations $\Box \Box$ The investigated molecule under consideration possesses CH₃ groups in third position of MNA in the ring. For the assignments of CH₃ group frequencies one can expected that nine fundamentals can be associated to each CH₃ group, namely three stretching, three bending, two rocking modes and a single torsional mode describe the motion of methyl group. The CH₃ symmetric stretching frequency is identified at 2983 cm⁻¹ in the FTIR and 2951 cm⁻¹ in the FT-Raman spectrum for MNA. The CH₃ inplane stretching vibrations are identified at 2929 cm⁻¹ in FT-Raman spectrum and 2940 cm⁻¹ in the FTIR spectrum for MNA. The CH₃ symmetric bending and CH₃ in-plane bending frequencies are attributed at 1274 cm⁻¹ and 1462, 1441 cm⁻¹ in the FTIR spectrum and 1263 cm⁻¹ FT-Raman spectrum for MNA. These assignments are supported by literature ²². The in-plane rocking and out-of-plane rocking modes of CH3 group are found at 916 cm⁻¹ and 1154 cm⁻¹ in the FTIR spectrum and 883, 1164 cm⁻in the FT-Raman spectrum for MNA. The bands obtained at 2894, 2846 cm^{-1} and 1130 cm^{-1}

in the FTIR and 2847 cm⁻¹ and 1142 cm⁻¹ FT-Raman spectrum for MNA assigned to CH_3 out-of-plane stretching and CH_3 out-of-plane bending modes, respectively. The assignment of the bands at 215, 176 cm⁻¹ FT-Raman spectrum for MNA attributed to methyl twisting mode.

*NO*₂ vibrations \Box There are six normal modes due to an NO₂ group namely asymmetric NO₂ stretching, symmetric NO₂ stretching, NO₂ rocking, NO₂ wagging and NO₂ torsion. The NO₂ asymmetric stretching vibration band range is 1625 - 1540 cm⁻¹ and that of symmetric stretching vibration is 1400 - 1360 cm⁻¹²². In the present investigation, the NO₂ asymmetric stretching vibrations for both IR and Raman spectrum is observed at 1613 cm⁻¹. The symmetric stretching vibration is observed at 1285 cm⁻¹ in FTIR spectrum for MNA. The NO₂ scissoring vibration is observed at 851 cm⁻¹ in IR spectrum. The deformation vibrations of NO₂ group (rocking, wagging and twisting) contribute to several normal modes in the low frequency region. These bands are also found well within the characteristic region and are summarized in Table 4.

5.1. Molecular properties

The thermodynamic parameters namely heat capacity, entropy, rotational constants, dipole moments, vibrational and zero-point vibrational energies of the compound have also been computed at ab initio HF and DFT/B3LYP levels using 6-311++G(d,p) as basis set and are presented in Table 5 for MNA, respectively. The difference in the values calculated by both the methods is marginal. The variation in the ZPVE seems to be insignificant. The ZPVE is much lower by the DFT/B3LYP method than by the HF method. The total energy and the change in the total entropy of the molecule at room temperature are also presented. Dipole moment reflects the molecular charge distribution and is given as a vector in three dimensions. Therefore, it can be used as descriptor to depict the charge movement across the molecule. Direction of the dipole moment vector in a molecule depends on the centers of positive and negative charges. Dipole moments are strictly determined for neutral molecules. For charged systems, its value depends on the choice of origin and molecular orientation. As a result of HF and DFT (B3LYP) calculations, the highest dipole moment is observed for B3LYP/6-311++G(d,p) level whereas the lowest one is observed for HF/6-311++G(d,p) level in the molecule.

The total dipole moment of MNA determined by HF and B3LYP level using 6-311++G(d,p) basis set is 0.9606, 1.2182 Debye, respectively.

6. Prediction of First Hyperpolarizability

The first hyperpolarizability (β_0) of this novel molecular system and the related properties (β_0 , α_0) of MNA are calculated using the B3LYP/ 6-311++G(d,p) basis set, based on the finite field approach. In the presence of an applied electric field, the energy of a system is a function of the electric field. The first hyperpolarizability is a third-rank tensor that can be described by a 3 × 3 × 3 matrix²³. The 27 components of the 3D matrix can be reduced to 10 components due to the Kleinman symmetry. It can be given in the lower tetrahedral. The components of β are defined as the coefficients in the Taylor series expansion of the energy in the external electric field.

When the external electric field is weak and homogenous, this expansion becomes:

$$\mathbf{E} = \mathbf{E}_{0} - \boldsymbol{\mu}_{\alpha} \mathbf{F}_{\alpha} - \frac{1}{2} \boldsymbol{\alpha}_{\alpha\beta} \mathbf{F}_{\alpha} \mathbf{F}_{\beta} - \frac{1}{6} \boldsymbol{\beta}_{\alpha} \boldsymbol{\beta}_{\gamma} \mathbf{F}_{\alpha} \mathbf{F}_{\beta} \mathbf{F}_{\gamma} + \dots$$

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Table 4. The observed (FTIR and FT-Raman) and calculated (unscaled and scaled) fundamental harmonic frequencies (cm⁻¹), force constant (mdyn A⁻¹), infrared - intensity (km/mol), Raman activity (Å amu⁻¹) and probable assignments of 3-methyl-2-nitroanisole are analysed based on SQM force field calculation using HF/6-311++G(d,p) and B3LYP/6-311++G(d,p) method and basis set calculations.

| Observed free | juencies | HF/6-311+ | ++G(d , p) | | | | B3LYP/6-311++G(d,p) | | | | Assignments (TED %) | |
|---------------|----------|-----------|----------------------------|-------------|--------------|----------------|---------------------|--------|-------------|--------------|---------------------|------------------------|
| FTIR | FT-Raman | Unscaled | Scaled | Force Const | IR intensity | Raman activity | Unscaled | Scaled | Force Const | IR intensity | Raman activity | |
| - | 3101w | 3382 | 3221 | 7.3969 | 6.8601 | 165.695 | 3213 | 3105 | 6.6446 | 3.4982 | 118.02 | vCH(99) |
| - | 3082w | 3370 | 3258 | 7.3202 | 14.0719 | 14.0719 | 3212 | 3088 | 6.5533 | 12.3767 | 154.92 | vCH(98) |
| - | 3030vw | 3353 | 3222 | 7.2153 | 3.2637 | 57.4249 | 3192 | 3035 | 6.4468 | 3.1175 | 60.322 | υCH(97) |
| 2983vw | - | 3321 | 3196 | 7.1914 | 27.3335 | 100.915 | 3146 | 2987 | 6.4142 | 12.9628 | 102.83 | CH3 ss(90), CH3ips(10) |
| - | 2951w | 3284 | 3129 | 7.0208 | 15.8950 | 58.8190 | 3129 | 2954 | 6.3187 | 12.8697 | 56.891 | CH3ss(87), CH3ips(13) |
| 2940vw | - | 3281 | 3100 | 7.0027 | 26.9387 | 42.8831 | 3094 | 2946 | 6.0241 | 6.3240 | 62.982 | CH3ips(79), CH3 ss(20) |
| - | 2929vw | 3270 | 3087 | 6.9478 | 12.1930 | 59.1378 | 3077 | 2935 | 6.1767 | 29.5562 | 56.985 | CH3ips(71), CH3 ss(26) |
| 2894vw | - | 3203 | 3058 | 6.2685 | 16.9173 | 173.631 | 3038 | 2899 | 5.6411 | 12.8974 | 204.01 | CH3ops(74), vCC(24) |
| 2846vw | 2847w | 3201 | 3074 | 6.2460 | 58.0375 | 84.1631 | 3054 | 2858 | 5.5360 | 49.9870 | 158.51 | CH3ops(79), vCC(20) |
| 1613ms | 1613w | 1827 | 1766 | 22.796 | 624.865 | 3.4224 | 1696 | 1619 | 12.2869 | 92.9804 | 10.296 | NO2ass(89), CH3ops(11) |
| 1586ms | - | 1789 | 1682 | 14.563 | 103.540 | 16.6865 | 1669 | 1597 | 8.4660 | 71.0762 | 32.102 | vCC(63), CH3ops(32) |
| - | 1559vw | 1780 | 1595 | 11.140 | 67.6607 | 18.7550 | 1595 | 1565 | 18.4206 | 275.226 | 12.282 | vCC(88), NO2ass(12) |
| - | 1538w | 1756 | 1575 | 18.392 | 107.991 | 20.7632 | 1565 | 1544 | 2.2006 | 69.8124 | 7.5901 | υCC(77), CH3ipb(21) |
| 1522vs | - | 1699 | 1554 | 3.3958 | 82.3702 | 3.0327 | 1556 | 1528 | 1.8615 | 35.3634 | 4.9048 | vCC(73), CH3ipb(22) |
| - | 1489vs | 1671 | 1543 | 1.6777 | 9.7713 | 7.4470 | 1510 | 1498 | 1.6404 | 14.1779 | 1.0648 | vCC(85) |
| 1462s | - | 1666 | 1523 | 2.2926 | 16.6797 | 3.9792 | 1501 | 1468 | 1.3735 | 8.5311 | 12.606 | CH3ipb(80), vCC(19) |
| 1441s | - | 1650 | 1593 | 1.6288 | 8.4194 | 11.2090 | 1545 | 1446 | 1.5087 | 24.5830 | 5.8904 | CH3ipb(82), vCC(18) |
| 1374ms | 1374w | 1634 | 1588 | 1.8083 | 17.7256 | 6.8157 | 1497 | 1385 | 1.6906 | 6.7976 | 2.9291 | vCC(79), NO2ss(20) |
| 1310vs | 1311ms | 1626 | 1578 | 1.9015 | 11.2427 | 4.0213 | 1461 | 1324 | 2.1681 | 22.2932 | 1.4018 | vCC(73), CH3sb(19) |
| 1285s | - | 1612 | 1450 | 2.5002 | 16.4548 | 3.3615 | 1420 | 1288 | 1.4705 | 1.4454 | 14.042 | NO2ss(75), vCC(25) |
| 1274vw | - | 1552 | 1437 | 1.7368 | 0.6752 | 6.8007 | 1407 | 1276 | 14.6833 | 143.327 | 22.193 | CH3sb(64), vCN(30) |
| - | 1263w | 1429 | 1407 | 4.0781 | 141.329 | 19.3044 | 1339 | 1266 | 5.7839 | 60.3669 | 14.314 | CH3sb(65), vCO(32) |
| 1255s | - | 1388 | 1386 | 1.9649 | 2.7693 | 2.0602 | 1308 | 1258 | 3.7019 | 150.593 | 19.388 | vCN(66), CH3sb(36) |
| 1189vw | - | 1324 | 1297 | 1.6677 | 1.5622 | 3.0395 | 1278 | 1193 | 1.8984 | 2.9039 | 3.1854 | υCO(72), CH3opr(22) |
| 1178vw | - | 1306 | 1260 | 1.4878 | 4.6140 | 1.1436 | 1208 | 1182 | 1.1473 | 1.2083 | 9.4959 | υCO(78), CH3opr(21) |
| - | 1164vs | 1279 | 1216 | 1.2068 | 2.5699 | 3.0848 | 1196 | 1166 | 1.0774 | 0.0425 | 1.0606 | CH3opr(74), vCO(20) |
| 1154vs | - | 1255 | 1201 | 2.9981 | 10.5563 | 1.5368 | 1168 | 1157 | 1.0201 | 0.9366 | 2.2797 | CH3opr(73), CH3opb(22) |
| - | 1142w | 1232 | 1194 | 2.0080 | 5.0160 | 6.8829 | 1147 | 1145 | 1.6734 | 7.2738 | 1.0146 | CH3opb(77), Rasym(26) |
| 1130vs | - | 1188 | 1168 | 3.1425 | 127.927 | 8.6351 | 1118 | 1135 | 3.0546 | 97.978 | 1.6004 | CH3opb(73), Rsymd(21) |
| 1091ms | - | 1166 | 1144 | 1.5915 | 5.9780 | 5.1572 | 1102 | 1093 | 2.5356 | 12.444 | 24.322 | Rasym(72), Rtrigd(22) |
| 1070ms | - | 1162 | 1111 | 1.6635 | 19.2732 | 6.1602 | 1099 | 1077 | 1.0122 | 2.1580 | 0.6462 | Rsymd(71), bCH(29) |

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| 4332 | | | | K.Pa | rivathini et al./ | Elixir Comput | ational Phys | ics 102 (2) | 017) 44325-4 | 4336 | | |
|-------|--------|------|------|--------|-------------------|---------------|--------------|-------------|--------------|---------|--------|--------------------------|
| - | 1024vs | 1116 | 1076 | 1.0191 | 3.4024 | 0.8975 | 1066 | 1029 | 1.1639 | 17.693 | 5.6861 | Rtrigd(70), bCH(28) |
| - | 1001w | 1106 | 1069 | 1.4157 | 12.3866 | 10.5245 | 1045 | 1009 | 0.7231 | 0.2010 | 0.0488 | bCH(61), bCH(35) |
| 998ms | - | 1069 | 1039 | 1.4715 | 7.6208 | 0.7001 | 1023 | 1007 | 1.9245 | 4.8174 | 0.7472 | bCH(68), CH3ipr(29) |
| 976vs | - | 1054 | 1023 | 1.0267 | 3.3566 | 1.6017 | 998 | 987 | 0.6272 | 0.3401 | 0.2340 | bCH(67), CH3ipr(30) |
| 916vs | - | 1023 | 1001 | 5.1768 | 31.5553 | 7.9995 | 968 | 924 | 4.1843 | 29.261 | 7.0928 | CH3ipr(66), bCN(32) |
| - | 883vw | 903 | 896 | 1.6557 | 23.8658 | 0.4911 | 914 | 887 | 1.5333 | 18.2664 | 0.4762 | CH3ipr(65), NO2sciss(30) |
| 869ms | - | 887 | 866 | 0.8183 | 35.0342 | 0.9669 | 901 | 873 | 0.5880 | 35.442 | 0.4791 | bCN(62), NO2wag(31) |
| 851vs | - | 872 | 861 | 1.7372 | 14.7314 | 3.6330 | 988 | 855 | 2.1935 | 7.8292 | 9.1458 | NO2sciss(61), ωCH(22) |
| 840ms | - | 860 | 851 | 2.1050 | 18.5362 | 4.8287 | 973 | 844 | 1.3043 | 11.5592 | 0.2433 | NO2wag(60), ωCH(27) |
| 776ms | - | 830 | 835 | 1.1987 | 13.6487 | 5.3267 | 934 | 778 | 1.2111 | 15.9929 | 6.3851 | ωCH(51), bCC(35) |
| - | 752vw | 812 | 809 | 0.9921 | 0.8069 | 1.0931 | 915 | 756 | 0.9718 | 2.2649 | 2.5479 | ωCH(58), bCO(32) |
| 745vs | - | 805 | 800 | 1.1038 | 8.6232 | 13.5470 | 889 | 749 | 0.7108 | 0.5416 | 1.5742 | ωCH(57), bCO(39) |
| 739w | - | 790 | 785 | 1.0700 | 1.7676 | 4.7339 | 830 | 743 | 0.8465 | 0.2442 | 7.6360 | NO2rock(56), ωCH(38) |
| 728w | - | 777 | 771 | 0.6233 | 6.9910 | 1.5021 | 781 | 732 | 0.6188 | 1.1593 | 5.0525 | bCC(44), tRasym(32) |
| - | 680vw | 768 | 755 | 0.6397 | 6.4354 | 1.0886 | 719 | 684 | 0.4853 | 1.6928 | 0.4246 | bCO(51), tRsymd(38) |
| 640vw | - | 745 | 735 | 0.5584 | 2.8908 | 3.1299 | 677 | 644 | 0.4957 | 0.1962 | 3.6634 | bCO(50), tRtrigd(49) |
| 622vw | 622w | 723 | 712 | 0.2757 | 0.7238 | 1.9933 | 635 | 624 | 0.2361 | 1.8036 | 4.4869 | t Rasym(44), ωCC(41) |
| 598w | - | 708 | 698 | 0.2687 | 2.2776 | 0.3882 | 660 | 599 | 0.0817 | 0.9717 | 0.4309 | t Rsymd(45), ωCO(40) |
| - | 569w | 690 | 678 | 0.1066 | 1.0396 | 2.1909 | 631 | 573 | 0.1092 | 4.9423 | 0.0705 | t Rtrigd(50), ωCN(46) |
| - | 554w | 678 | 663 | 0.1519 | 3.4704 | 0.3266 | 607 | 559 | 0.1075 | 0.3251 | 2.4707 | ωCC(40) |
| - | 524w | 654 | 645 | 0.0834 | 0.8627 | 3.7644 | 601 | 527 | 0.0348 | 0.1472 | 1.9280 | ωCO(45) |
| - | 479vw | 634 | 623 | 0.0146 | 0.4515 | 0.1832 | 563 | 483 | 0.0972 | 2.0378 | 1.4083 | ωCO(47) |
| - | 385vw | 623 | 612 | 0.0160 | 1.3762 | 0.2214 | 525 | 389 | 0.0099 | 0.3523 | 0.1913 | ωCN(46) |
| - | 342vs | 455 | 412 | 0.0230 | 5.1544 | 0.3923 | 503 | 347 | 0.0328 | 0.0900 | 2.0186 | NO2twist(49) |
| - | 215s | 427 | 412 | 0.0153 | 4.5374 | 1.3341 | 444 | 218 | 0.0101 | 5.2598 | 0.7334 | tCH3(47) |
| - | 176ms | 408 | 399 | 0.0329 | 0.0709 | 3.2964 | 420 | 179 | 0.0157 | 0.5548 | 4.1768 | tCH3(43) |

Abbreviations: v - stretching; b - in-plane bending; ω - out-of-plane bending; asymd - asymmetric; symd - symmetric; t - torsion; trig - trigonal; w - weak; vw - very week; vs - very strong; s - strong; ms - medium strong; ss - symmetric stretching; ass - asymmetric stretching; ips - in-plane stretching; ops - out-of-plane stretching; sb - symmetric bending; ipr - in-plane rocking; opr - out-of-plane rocking; opb - out-of-plane bending.

Table 5. The thermodynamic parameters of 3-methyl-2nitroanisole calculated at HF/6-311++G(d,p) with and B3LYP/6-311++G(d,p) and 6-311++G(d,p) method and basis set calculations.

| Parameters | HF/ | B3LYP/ |
|---|-------------|-------------|
| | 6- | 6- |
| | 311++G(d,p) | 311++G(d,p) |
| Optimized global minimum | -587.1106 | -590.7489 |
| Energy, (Hartrees) | | |
| Total energy(thermal), E _{total} | 116.454 | 108.908 |
| (kcal mol ⁻¹) | | |
| Translational | 0.889 | 0.889 |
| Rotational | 0.889 | 0.889 |
| Vibrational | 28.965 | 107.130 |
| Molar capacity at constant | | |
| volume, | | |
| $(cal mol^{-1} k^{-1})$ | | |
| Total | 38.377 | 41.220 |
| Translational | 2.981 | 2.981 |
| Rotational | 2.981 | 2.981 |
| Vibrational | 34.124 | 35.259 |
| Entropy | | |
| Total | 102.089 | 105.246 |
| Translational | 41.248 | 41.258 |
| Rotational | 30.397 | 30.519 |
| Vibrational | 29.897 | 33.479 |
| Zero point vibrational energy, | 109.7325 | 101.7850 |
| (Kcal mol ⁻¹) | | |
| Rotational constants (GHZ) | | |
| Α | 1.1839 | 1.1822 |
| В | 1.0719 | 0.9992 |
| С | 0.6168 | 0.5857 |
| Rotational temperature (Kelvin) | 0.0568 | 0.0567 |
| _ | 0.0514 | 0.0479 |
| | 0.0296 | 0.0281 |

where E_0 is the energy of the unperturbed molecules, F_{α} the field at the origin and μ_{α} , $\alpha_{\alpha\beta}$ and $\beta_{\alpha}\beta_{\gamma}$ are the components of dipole moment, polarizability and the first hyperpolarizabilities, respectively. The total static dipole moment μ , the mean polarizability α_0 and the mean first hyperpolarizability β_0 , using the *x*, *y*, *z* components they are defined as:

$$\alpha_{0} = \frac{\left(\alpha_{xx} + \alpha_{yy} + \alpha_{zz}\right)}{3}$$

$$\alpha = 2^{-\frac{1}{2}} \left[\left(\alpha_{xx} - \alpha_{yy}\right)^{2} + \left(\alpha_{yy} - \alpha_{zz}\right)^{2} + \left(\alpha_{zz} - \alpha_{xx}\right)^{2} + 6\alpha_{xx}^{2} \right]^{\frac{1}{2}}$$

$$\beta_{0} = \left(\beta_{x}^{2} + \beta_{y}^{2} + \beta_{z}^{2}\right)^{\frac{1}{2}}$$

$$\beta_{vec} = \frac{3}{5\left[\left(\beta_x^2 + \beta_y^2 + \beta_z^2\right)^{\frac{1}{2}}\right]}$$

where

where

 $\begin{aligned} \beta_{x} &= \beta_{xxx} + \beta_{xyy} + \beta_{xzz} \\ \beta_{y} &= \beta_{yyy} + \beta_{yxx} + \beta_{yzz} \\ \beta_{z} &= \beta_{zzz} + \beta_{zxx} + \beta_{zyy} \end{aligned}$

The β_0 components of GAUSSIAN 09W output are reported in atomic units and therefore the calculated values are converted into e.s.u. units (1 a.u. = 8.3693×10^{-33} e.s.u.). The calculated value of hyperpolarizability and polarizability of MNA are tabulated in Table 6, respectively.

The title molecules are an attractive object for future studies of non-linear optical properties.

| Table 6. Nonlinear optical properties of 3,4-di | methylanisole |
|---|---------------|
| and 3-methyl-2-nitroanisole calculated at HF/6 | 6-311++G(d,p) |
| with B3LYP/6-311+G(d,p) and 6-311++G(d,p) | method and |
| basis set calculations. | |

| NLO behaviour | MNA | | | | | |
|---------------------------------------|--------------------------|--------------------------|--|--|--|--|
| | HF/ | B3LYP/ | | | | |
| | 6-311+G(d,p) | 6- | | | | |
| | | 311++G(d,p) | | | | |
| Dipole moment(µ) | 0.9606Debye | 1.2182Debye | | | | |
| Mean polarizability (α) | 0.8846×10^{-30} | 0.9909 × 10 ⁻ | | | | |
| | esu | ³⁰ esu | | | | |
| Anisotropy of the | 1.7675×10^{-30} | 1.2708×10^{-1} | | | | |
| polarizabilty (Δ_{α}) | esu | ³⁰ esu | | | | |
| First hyperpolarizability (β) | 1.6177×10^{-30} | 2.9059×10^{-10} | | | | |
| | esu | ³⁰ esu | | | | |
| Vector – first | 0.9706×10^{-30} | 1.7435 × 10 ⁻ | | | | |
| hyperpolarizability (β_{vec}) | esu | ³⁰ esu | | | | |

7. HOMO-LUMO Band Gap

This electronic absorption corresponds to the transition from the ground to the first excited state and is mainly described by one electron excitation from the highest occupied molecular orbital (HOMO) to the lowest unoccupied molecular orbital (LUMO)24. Many organic molecules, containing conjugated π electrons are characterized by large values of molecular first hyper polarizabilities and analyzed by means of vibrational spectroscopy. In most of the cases, even in the absence of inversion symmetry, the band in the Raman spectrum is weak in the IR spectrum and vice-versa. But the intra molecular charge from the donor to acceptor group through a single-double bond conjugated path can induce large variations of both the molecular dipole moment and the molecular polarizability, making IR and Raman activity strong at the same time. The analysis of the wave function indicates that the electron absorption corresponds to the transition from the ground to the first excited state and is mainly described by one-electron excitation from the highest occupied molecular orbital (HOMO) to the lowest unoccupied orbital (LUMO). In MNA, the HOMO is located over heterocyclic ring and the HOMO-LUMO transition implies an electron density transfer to the CH₃, NO₂ group from heterocyclic ring and oxygen atom, whereas in MNA the HOMO is located over benzene ring, especially on nitro and oxygen atom, and the HOMO-LUMO transition implies an electron density transfer to the heterocyclic ring from nitro group and oxygen atom. Moreover, the composition of HOMO and LUMO for MNA is shown in Fig. 5, respectively. The HOMO-LUMO energy gap of MNA are calculated at B3LYP/6-311++G(d,p) level, which reveals that the energy gap reflects the chemical activity of the molecules. The LUMO as an electron acceptor (EA) represents the ability to obtain an electron (ED) and HOMO represents ability to donate an electron (ED). The ED groups to the efficient EA groups through π -conjugated path. The strong charge transfer interaction through π -conjugated bridge results in substantial ground state Donor-Acceptor (DA) mixing and the appearance of a charge transfer band in the electron absorption spectrum. The HOMO and LUMO energy gap explains the fact that

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eventual charge transfer interaction is taking place within the title molecules.



Fig 5. HOMO-LUMO plot of 3-methyl-2-nitro anisole.

7.1. Global and local reactivity descriptors

Based on density functional descriptors global chemical reactivity descriptors of compounds such as hardness, chemical potential, softness, electronegativity and electrophilicity index as well as local reactivity have been defined ²⁵. Pauling introduced the concept of electronegativity as the power of an atom in a compound to attract electrons to it. Hardness (η), chemical potential (μ) and electronegativity (χ) and softness are defined follows.

$$\begin{split} \eta &= \frac{1}{2} \left(\partial^2 E / \partial N^2 \right)_{V(r)} = \frac{1}{2} \left(\partial \mu / \partial N \right)_{V(r)} \\ \mu &= \left(\partial E / \partial N \right)_{V(r)} \\ \chi &= -\mu = - \left(\partial E / \partial N \right)_{V(r)} \end{split}$$

where E and V(r) are electronic energy and external potential of an N-electron system respectively. Softness is a property of compound that measures the extent of chemical reactivity. It is the reciprocal of hardness.

$$S = \frac{1}{\eta}$$

Using Koopman's theorem for closed-shell compounds, $\eta,$ μ and χ can be defined as

$$\eta = \frac{(I-A)}{2}$$
$$\mu = \frac{-(I+A)}{2}$$
$$\chi = \frac{(I+A)}{2}$$

where A and I are the ionization potential and electron affinity of the compounds respectively. Electron affinity refers to the capability of a ligand to accept precisely one electron from a donor. However in many kinds of bonding viz. covalent hydrogen bonding, partial charge transfer takes places. Recently Parr *et al.* ²⁶ have defined a new descriptor to quantity the global electrophilic power of the compound as electrophilicity index ($_{(\Omega)}$), which defines a quantitative classification of the global electrophilic nature of a compound have proposed electrophilicity index ($_{(\Omega)}$) as a measure of energy lowering due to maximal electron flow between donor and acceptor. They defined electrophilicity index ($_{(\Omega)}$) as follows.

$$\omega = \mu^2/2\eta$$

The usefulness of this new reactivity quantity has been recently demonstrated in understanding the toxicity of various pollutants in terms of their reactivity and site selectivity. The calculated value of electrophilicity index describes the biological activity for MNA respectively. All the calculated values of HOMO-LUMO, energy gap, ionization potential, Electron affinity, hardness, potential, softness and electrophilicity index are shown in Table 7, respectively.

8. Nbo Analysis

NBO analysis gives information about interactions in both filled and virtual orbital spaces that could enhances the analysis of intra- and intermolecular interactions. The larger the $E^{(2)}$ (energy of hyperconjugative interactions) value, the more intensive is the interaction between electron donors and electron acceptors, i.e. the more donating tendency from electron donors to electron acceptors the greater the extent of conjugation of the whole system. Delocalization of electron density between occupied Lewis-type (bond or lone pair) NBO orbitals and formally unoccupied (anti-bond or Rydberg) non-Lewis NBO orbitals correspond to a stabilizing donor-acceptor interaction. NBO analysis has been performed on the title molecule at the DFT level in order to elucidate the intra-molecular, rehybridization and delocalization of electron density within the molecule. The larger $E^{(2)}$ value, the more intensive is the interaction between electron donors and acceptors, i.e., the more donation tendency from electron donors to electron acceptors and the greater the extent of conjugation of the whole system. DFT (B3LYP/6-311++G(d,p)) level computation is used to investigate the various second-order interactions between the filled orbitals of one subsystem and vacant orbitals of another subsystem, which is a measure of the delocalization or hyperconjugation²⁷. NBOs are localized electron pair orbitals for bonding pairs and lone pairs. The hybridization of the atoms and the weight of each atom in each localized electron pair bond are calculated in the idealized Lewis structure. A normal Lewis structure would not leave any antibonding orbitals, so the presence of antibonding orbitals shows deviations from normal Lewis structures. Anti bonding localized orbitals are called non-Lewis NBOs. In order to study the small deviations from idealized Lewis structure, the Donor-Acceptor interaction approach is adopted. In MNA, $\pi(C3-C4) \rightarrow \pi^*(C1-C2)$, $\pi(C5-C6) \rightarrow \pi^*(C1-C2)$ interaction is seen to give a strong stabilization 41.84 and 52.46 kJ/mol. This strong stabilization denotes the larger delocalization. The interesting interactions in MNA molecule are LP1N12, LP2O7with that of antibonding O13-O14, C1-C6. These two interactions result the stabilization energy of 162.54, 13.84 kJ/mol respectively. This highest interaction around the ring can induce the large bioactivity in the molecule. This shows that the lone pair orbital participates in electron donation in the molecule. The calculated values of $E^{(2)}$ are shown in Table 8 for MNA, respectively.

9. ¹³C AND ¹H NMR SPECTRAL ANALYSIS

The molecular structure of MNA is optimized by using B3LYP method with 6-31++G basis set. Then, GIAO ¹³C calculations of the title compound are calculated and compared with experimental values ²⁸ are shown in Table 9. Relative chemical shifts are then estimated by using the corresponding TMS shielding calculated in advance at the theoretical level as reference. Changes in energy needed to flip protons are called chemical shifts. The location of chemical shifts (peaks) on a NMR spectrum are measured from a reference point that the hydrogen in a standard reference compound -(CH₃)₄Si or tetramethylsilane (TMS)- produce. The amount of energy necessary to flip protons in TMS is assigned the arbitrary value of zero δ . Chemical shifts are measured in parts per million magnetic field strength difference (δ -scale), relative to TMS. The experimental values of MNA for ¹H and ¹³C isotropic chemical shielding for TMS are 161.30, 163.80 ppm, respectively ²⁸. All the calculations are performed using Gauss view molecular visualization program and Gaussian 09W program package. The result shows that the range ¹³C NMR chemical shift of the typical organic compound usually is > 100 ppm, the accuracy ensures reliable interpretation of spectroscopic parameters. It is true from the above literature value in our present study, that the title compound also shows the same. In practice, it is easier to fix the radio wave frequency and vary the applied magnetic field than it is to vary the radio wave frequency. The magnetic field "felt" by a hydrogen atom is composed of both applied and induced fields. The induced field is a field created by the electrons in the bond to the hydrogen and the electrons in nearby π bonds. When the two fields reinforce each other, a smaller applied field is required to flip the proton. In this situation, a proton is said to be deshielded. When the applied and induced fields oppose each other, a stronger field must be applied to flip the proton. In this state, the proton is shielded. Electronegative atoms such as CH₃, O, NO2 and halogens deshield hydrogen. The extent of deshielding is proportional to the electronegativity of the heteroatom and its proximity to the hydrogen. These nitro and oxygen atoms show electronegative property, so that the chemical shift of C1, C2, C3, C4, C5, C6 for MNA seems to be 161.28, 163.80, 141.74, 121.87, 131.91 and 122.65 ppm. The chemical shift of C9 is greater than the other carbon values. This increase in chemical shift is due to the substitution of more electronegative oxygen and nitro atoms in the benzene ring. The presence of electronegative atom attracts all electron clouds of carbon atoms towards the oxygen and nitro atoms, which leads to deshielding of carbon atom and net result in increase in chemical shift value. The NMR shielding surfaces of C15, C8 is shown in this work the chemical shift (δ) for carbon atoms presented in the MNA in gas phase has been studied and theoretical ¹³C, ¹H-NMR isotropic shielding of carbon and Hydrogen atom. In the NMR shielding surfaces, the blue region represents shielding and red region represents de-shielding are shown in Fig. 6 for MNA, respectively.delocalization of electron density within the molecule.



Fig 6. NMR shielding surface of 3-methyl-2-nitro anisole.

10. Conclusion

The molecular structural parameters, thermodynamic properties and fundamental vibrational frequencies of the optimized geometry of 3-methyl-2-nitroanisole have been obtained from ab initio HF and DFT calculations. The theoretical results are compared with the experimental vibrations. The computed geometrical parameters are in good agreement with the observed X-ray diffraction data of similar compound. Although both types of calculations are useful to explain vibrational spectra of 3-methyl-2-nitroanisole, ab initio calculations at HF/6-311++G(d,p) level is found little poorer than DFT-B3LYP/6-311++G(d,p) level calculations. On the basis of agreement between the calculated and experimental results, assignments of all the fundamental vibrational modes of 3methyl-2-nitroanisole have been made for the first time in this investigation. The TED calculation regarding the normal modes of vibration provides a strong support for the frequency assignment. Therefore, the assignments proposed at higher level of theory with higher basis set with only reasonable deviations from the experimental values seem to be correct. HOMO and LUMO energy gap explains the eventual charge transfer interactions taking place within the molecule. Furthermore, the nonlinear optical, first-order hyperpolarizabilities and total dipole moment properties of the molecule show that the title molecule is an attractive object for future studies of nonlinear optical properties. NMR, NBO analysis have been performed in order to elucidate charge transfers or conjugative interaction, the intra-molecule rehybridization and delocalization of electron density within the molecule.

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